

# RESSALVA

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Silver (Ag) nanoclusters and lanthanides ( $\text{Er}^{3+}$ ,  $\text{Pr}^{3+}$ ,  $\text{Yb}^{3+}$ ) ions  
functionalized glassy materials for light sources, luminescent  
layers for silicon solar cell and waveguide application

TARCIO DE CASTRO SILVA

PhD Thesis

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## CERTIFICADO DE APROVAÇÃO

TÍTULO DA TESE: "Silver (Ag) nanoclusters and lanthanides ( $\text{Er}^{3+}$ ,  $\text{Pr}^{3+}$ ,  $\text{Yb}^{3+}$ ) ions functionalized glassy materials for light sources, luminescent layers for silicon solar cell and waveguide application"

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## **POSTER AND ORAL PRESENTATION IN NATIONAL AND INTERNATIONAL EVENTS**

Fares, H.; Silva, T. C.; Orives, J. R.; Franco, D. F.; Nalim, M. White light generation from Ag nanoclusters doped oxyfluoride glass for broadband light source. 10<sup>th</sup> International Conference on Nanophotonics. 2017. (Poster – Best poster award).

Silva, T. C.; Manzani, Danilo; Ribeiro, S. J. L. Visible to near-infrared upconversion emission of highly content Er<sup>3+</sup>-doped fluoroindate glasses and glass/polymer composites for solar cells applications. Advanced School on Glass and Glass-Ceramics. 2016; (Poster).

Silva, T. C.; Manzani, D.; Ribeiro, S. J. L. XXXVIII Encontro Nacional de Física da Matéria Condensada. Near Infrared upconversion luminescence in RE (RE=Er<sup>3+</sup>, Yb<sup>3+</sup>) ions doped fluoroindate glasses and YVO<sub>4</sub> nanoparticles. 2015; (Poster).

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## **DEDICATION**

I would like to express my eternal thoughts and love for those who passed away and made part in my life. Your faces are carved in my memories father "Alcebiades", uncle "Deilane" and grandmother "Caroline". There is no one single day that you are not remembered.

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## RESUMO

Materiais vítreos se tornaram um dos mais atrativos materiais devido a fácil e rápida preparação. Durante as últimas décadas, várias aplicações têm sido exploradas e a viabilidade para mudança de frequência, conversão ascendente (upconversion) e descendente de energia (downconversion), geração de radiação visível com similaridade das emissões cromáticas azul, verde e vermelho; guia de onda, camadas luminescentes para células solares, geração de segundo harmônico e outros, estão publicadas. Propriedades estruturais como baixa energia de vibração estrutural (fônon), estabilidade química e ampla transparência óptica, tem sido o principal alvo para explorar uma excelente matriz para a dopagem com terras-raras. Por exemplo, íons  $\text{Er}^{3+}$  é um dos candidatos mais adequados para a propriedade de “upconversion” devido a absorção na região do infravermelho em 0,9 e 1,5  $\mu\text{m}$  e consequente emissão na região do visível e infravermelho. Devido às suas propriedades espectroscópicas, os lantanídeos, geralmente, exibem emissões com linhas estreitas devido à fraca interação com o campo cristalino da matriz. Por outro lado, íons metálicos como a prata (Ag), exibem uma ampla banda de emissão e excitação. A propriedade mais interessante dos íons prata é que a emissão visível pode ser ajustada de acordo com a forma e tamanho das espécies de pratas presentes em uma amostra. Esta propriedade abre um largo espectro de aplicações nas quais a radiação visível pode ser explorada. Neste sentido, um dos principais objetivos dessa tese é mostrar a preparação de materiais vítreos contendo íons lantanídeos como  $\text{Er}^{3+}$  e  $\text{Pr}^{3+}/\text{Yb}^{3+}$  e /ou íons prata e nanoclusters de prata para aplicação em camadas luminescentes para células solares, geração de luz branca, guia de onda e investigação dos processos de “upconversion” e “downconversion”. Para este propósito, materiais vítreos com composição molar  $40\text{InF}_3\text{-}20\text{ZnF}_2\text{-}20\text{Sr}\text{-}20\text{BaF}_2$  (vidros fluoroindatos) e  $50\text{NaPO}_3\text{-}25\text{ZnF}_2\text{-}15\text{CdF}_2\text{-}10\text{YF}_3$  (vidros fluorofosfatos) foram selecionados. Os vidros fluoroindatos foram dopados com íons  $\text{Er}^{3+}$  com diferentes concentrações e os vidros fluorofosfatos com diferentes concentrações de pratas. Composições com diferentes concentrações de prata e os lantanídeos  $\text{Pr}^{3+}/\text{Yb}^{3+}$  também foram preparadas. As amostras vítreas foram preparadas pelo método de fusão-choque térmico. Os materiais de partida foram homogeneizados em almofariz. A fusão dos compostos para os vidros fluoroindatos foram realizadas em tubo de platina a 850 °C durante 20

minutos e vertidos em um molde metálico pré-aquecido em 260 °C, enquanto que a obtenção dos vidros fluorofosfatos se deu com fusão a 1000 °C durante 30 minutos e vertidos em molde pré-aquecido a 300 °C. A primeira parte desta tese mostra a preparação e avaliação das propriedades estruturais e espectroscópicas dos materiais vítreos. Absorção, excitação e espectros de emissão, o tamanho e morfologia dos nanoclusters de prata (para os vidros fluorofosfatos) bem como as propriedades térmicas foram analisadas. A segunda parte desta tese está relacionada com a aplicação dos vidros fluoroindatos como camadas luminescentes para o aumento da fotocorrente em uma célula solar de silício monocristalino e monofacial. Os vidros obtidos foram cortados e opticamente polidos de modo a obter amostras com as dimensões de 5.4 x 4.3 x 3.2 mm. As amostras foram colocadas na superfície de uma célula solar. Neste ponto, as diferentes concentrações de íons  $\text{Er}^{3+}$ , que estão associadas à diferentes intensidades de emissão, foram analisadas nas medidas de fotocorrente.

A terceira parte desta tese traz uma breve introdução aos nanoclusters de prata, que será a base para a compreensão dos resultados obtidos dos vidros fluorofosfatos. A maior parte dos resultados foram obtidos em colaboração com o “Institut de Chimie de la Matière Condensée de Bordeaux”, exceto os resultados das amostras preparadas para obtenção de radiação branca, capítulo 4. As amostras previamente preparadas com composição molar  $45\text{NaPO}_3\text{-}25\text{MgF}_2\text{-}15\text{CdF}_2\text{-}15\text{YF}_3\text{-}5\text{AgNO}_3$ , foram cortadas e polidas de modo a obter amostras com dimensão 7 x 4 x 2 mm. A estabilização dos nanoclusters de prata se deu através da técnica “Direct Laser Writing (DLW)”, utilizando laser de femtosegundo (fs)  $\text{KGW:Yb}$ , operando com linha espectral em 1030 nm e bombardeando a amostra com diferentes pulsos de energia. A irradiância laser foi controlada um modulador acústico-óptico permitindo a acumulação de  $N=10^5\text{-}10^6$  com energia variando entre 20 e 150 nJ. As microestruturas de prata gravadas nas amostras foram avaliadas de acordo com suas propriedades de absorção e emissão. Em adição, neste amostra foi avaliada a viabilidade como guia de onda. Por último, na amostra com composição molar  $(30\text{-}x)\text{NaPO}_3\text{-}30\text{MgF}_2\text{-}30\text{BaF}_2\text{-}10\text{YF}_3\text{-}0.5\text{Pr}\text{-}0.5\text{Yb}\text{-}x\text{Ag}$  ( $x = 0, 2, 4$  and  $10$ ) foi analisada a transferência de energia dos nanoclusters de prata e os terras-raras  $\text{Pr}^{3+}$  e  $\text{Yb}^{3+}$ . Os resultados apresentados nesta tese representam uma direção promissora para explorar materiais emissores com ampla emissão espectral com potencial uso em

conversão descendente de energia “downconversion”, camadas para aplicações fotovoltaicas, fósforos e guias de ondas.

Palavras-chave: Vidros fluoroindatos; vidros fluorofosfatos, “upconversion”, “downconversion”, transferência de energia, célula solar, nanoclusters de prata, “direct laser writing” e guias de onda.

## ABSTRACT

Glassy materials have become one of the most attractive materials owing their easy and fast preparation. During the last decades, several applications have been explored and the feasibility for frequency change, upconversion and downconversion, white light generation, waveguide, luminescent layers for solar cell, second harmonic generation and other, are reported. Structural properties like low energy phonon, chemical stability, broad transparency optical window have been the main target to explore a good host materials for lanthanides dopage. For example,  $\text{Er}^{3+}$  is one of the most suitable candidate for upconversion since it has 1,5  $\mu\text{m}$  and 0,9  $\mu\text{m}$  absorption. Owing to their spectroscopic properties, lanthanides, generally, show visible emissions with narrow lines due to the poor interaction with the crystalline field of the host material. On the other hand, metal ions such as silver; generally, shows a broadband emission and absorption. The most interesting spectroscopic properties of silver is that the visible emission may be tunable according to the size and shape of the silver species in the host materials. This property opens a large spectrum of application in which the visible light can be explored. In this sense, the main goal of this thesis is to show the preparation of glassy materials containing lanthanides such as  $\text{Er}^{3+}$ ,  $\text{Pr}^{3+}/\text{Yb}^{3+}$  and/or silver ions and nanoclusters for luminescent layers for solar cell and photonic applications. The glassy materials selected for such purposes were fluoroindate glasses and fluorophosphate glasses. The samples were prepared by the conventional melting-quenching method. The raw materials with high purity were weighted and mixed according to the following molar composition:  $40\text{InF}_3\text{-}20\text{ZnF}_2\text{-}20\text{Sr}\text{-}20\text{BaF}_2$  (fluoroindate glass) and  $50\text{NaPO}_3\text{-}25\text{ZnF}_2\text{-}15\text{CdF}_2\text{-}10\text{YF}_3$  (fluorophosphate glass). Besides, other compositions of fluorophosphate glasses were also synthesized. Fluoroindate glasses were doped with different  $\text{Er}^{3+}$  amount and the fluorophosphate glasses with different silver amounts. The raw materials of the fluoroindate glasses were melted in a platinum tube at 850 °C during 20 minutes and poured out in a mold pre-heated at 260 °C. On the other hand, the raw materials of the fluorophosphate glasses were melted in a platinum crucible at 1000 °C during 30 min and poured out in a mol pre-heated at 300 °C. The first part of this thesis was the preparation and the evaluation of the structural and spectroscopic properties of these glasses. Absorption, excitation and emission spectra, size and morphology of the silver nanoclusters (for

the fluorophosphate glasses) as well as thermal properties were analyzed. The second part of this thesis is related to the application of the fluorindate samples as luminescent layers for enhance photocurrent in a monocrystalline and monofacial silicon solar cell. The glasses obtained were cut and polished in order to obtain samples with dimensions 5.4 x 4.3 x 3.2 mm. The functional glass pieces were placed on the surface of the solar cell. At this point, the different  $\text{Er}^{3+}$  amount, which is associated with different upconversion intensity responses were evaluated in the photocurrent measures.

The third part of this thesis brings a brief introduction to the silver nanoclusters that will be the bases to understand the results obtained from the fluorophosphate glasses. Majority of these results was obtained in collaboration with "Institut de Chimie de la Matière Condensée de Bordeaux", except the results of the samples for white light generation, shown in the chapter 4. The samples previously prepared with molar composition  $45\text{NaPO}_3\text{-}25\text{MgF}_2\text{-}15\text{CdF}_2\text{-}15\text{YF}_3\text{-}5\text{AgNO}_3$ , were cut and polished in order to obtain samples with dimension 7 x 4 x 2 mm. The stabilization of silver nanoclusters was carried out by direct laser writing (DLW technique), using a femtosecond laser (KGW:Yb) operating at 1030 nm and delivering different energy pulses. The laser irradiance were controlled with an acousto-optic modulator enabling the accumulation of  $N=10^5\text{-}10^6$  pulses with energy pulses going from 20 nJ up to 150 nJ. The microstructures of silver nanoclusters were evaluated according their absorption and emission properties. Still, for this sample we also analyzed the feasibility as waveguiding material. For last but no least, the sample with molar composition  $(30\text{-}x)\text{NaPO}_3\text{-}30\text{MgF}_2\text{-}30\text{BaF}_2\text{-}10\text{YF}_3\text{-}0.5\text{Pr}\text{-}0.5\text{Yb}\text{-}x\text{Ag}$  ( $x = 0, 2, 4$  and  $10$ ) was prepared in order to evaluate the energy transfer among the silver species and the rare earths  $\text{Pr}^{3+}$  and  $\text{Yb}^{3+}$ . The results presented in this thesis represent a promising direction to explore broadband emitter's with potential use in down-conversion, layers for photovoltaic application, phosphors and waveguide. .

Keywords: Fluorindate glasses, fluorophosphate glasses, upconversion, downconversion, energy transfer, solar cell, silver nanoclusters, direct laser writing and waveguide.

*“Next time someone complains that you have made a mistake, tell him that may be a good thing. Because without imperfection, neither you nor I would exist.”*

*(Stephen Hawking, 2010)*

## LIST OF FIGURES

<b>Figure 1.1.</b> Schematic of a single-junction solar cell.....	26
<b>Figure 1.2.</b> Figure 1.2. Solar cell generation according to wafer-based and thin-film classification. (a) First generation: wafer-based; (b) second generation: commercial thin film; and (c) second generation: thin-film.....	28
<b>Figure 1.3.</b> Solar spectrum at surface hearth. The red region emphasize the portion of the terrestrial sunlight that is absorbed by a silicon SC with bandgap of 1.12 eV, and spectral regions that are lost by thermalization and below bandgap or transmitted photons.....	29
<b>Figure 1.4.</b> Losses in a single-junction SC. Process (a), thermalization photons and (b), transmitted photons.....	30
<b>Figure 1.5.</b> Electronic distribution of the Ce <sup>3+</sup> ion in a box diagram.....	33
<b>Figure 1.6.</b> Energy level for trivalent rare-earth ions (RE <sup>3+</sup> ) in LaF <sub>3</sub> crystal.....	34
<b>Figure 1.7.</b> Simplified energy levels for different two photons mechanisms in UC process. The excitation of the sensitizers ions are represented by orange arrows, non-radiative transition are represented by black dotted arrows and the purple arrows represent visible emission from the excited state E <sub>2</sub> to the ground state E <sub>0</sub> , respectively.....	36
<b>Figure 1.8.</b> Scheme of an ideal system for UC in a bifacial sola cell proposed by Trupke et al.....	37
<b>Figure 1.9.</b> Experimental scheme of an UC material and a solar cell proposed by Badescu. Adapted from the reference [BADESCU, 2008].....	38
<b>Figure 1.10.</b> (a) Experimental setup of a GaAs solar cell coupled with Yb <sup>3+</sup> /Er <sup>3+</sup> co-doped vitroceramic, and (b) energy levels showing red and green emission after energy transfer from Yb <sup>3+</sup> to Er <sup>3+</sup> . Figure (a) was adapted from the reference [GILBART, 1996].....	39
<b>Figure 1.11.</b> Energy levels of two Er <sup>3+</sup> ions showing UC process upon excitation at 1500 nm. Three subsequent ET from the sensitizer to the activator guarantee the population of the excited states in which UC emission is observed.....	40
<b>Figure 1.12.</b> Working principle of a luminescent layer placed at the top level of a silicon solar cell.....	41
<b>Figure 1.13.</b> Schematic diagram for DC in solar cell.....	42

<b>Figure 2.1.</b> Absorption spectrum of the sample IZSB-E7 (black line) and partial solar radiation spectrum AM1.5G (pink line). The dashed blue square emphasize the overlapping between Er <sup>3+</sup> absorption and solar spectrum emission over 1550 nm .....	52
<b>Figure 2.2.</b> Visible and near-infrared UC emission upon 1550 nm excitation and digital pictures of the samples (top level). The sample IZSB-E01 does not exhibit emission in the visible region and the UC emission is negligible. The intensity of the samples IZSB-E1, IZSB-E05 and IZSB-E01 was multiplied by a factor of 5.....	53
<b>Figure 2.3.</b> Schematic energy level diagram for Er <sup>3+</sup> -doped fluoroindate glass under excitation at 1550 nm .....	55
<b>Figure 2.4.</b> Decay time for green, red and near-infrared emission under excitation at 521 nm, figure (a), (b) and (c), respectively.....	56
<b>Figure 2.5.</b> Photocurrent analysis of the samples as a function of the power source excitation. Experimental setup for the photocurrent analysis (inset figure). The samples were place on the top of a Si solar cell and excited with a perpendicular beam at 1550 nm. All the measurements were performed in a dark room .....	57
<b>Figure 3.1.</b> Diagram energy levels of the Ag NC's proposed by Tikhomirov et al., [TIKHOMIROV, 2010]. The ground state (4d <sup>10</sup> 5s <sup>1</sup> ) and the first excited state (4d <sup>9</sup> 5s <sup>2</sup> ) may find another configuration in the literature according to the dominant nanocluster in a specific material.....	64
<b>Figure 3.2.</b> Electronic structure according to the different size of the silver metal. The electronic structure is broken down when the number of atoms increase in the metal.....	66
<b>Figure 3.3.</b> Excitation and emission spectra of oxyfluoride glass doped with 5% of AgNO <sub>3</sub> . Excitation are represented on the left side and emission on the right side as well.....	66
<b>Figure 3.4.</b> Ag NC's stabilized in poly(methacrylic)acid. Ag NC's are protected by the carboxyl groups. Photograph and emission spectra of the solutions.....	67
<b>Figure 3.5.</b> Fluorophosphate glasses containing different silver amounts. (a) Daylight photography of the samples and under UV irradiation, (b) and (c).....	68
<b>Figure 3.6.</b> (a) HRTEM of as-prepared oxyfluoride glass, (b) 6 hours and (c) 16 hours of heat treatment. (d) Size distribution of the Ag NC's. (e) Absorption spectra and (f) emission spectra excited at different wavelengths.....	70
<b>Figure 3.7.</b> Ag NC's and Ag NP's obtained in fluorophosphate glasses. (a) Photography of the samples in daylight and under ultraviolet radiation. (b) Transmission electron microscopy with 10 nm of resolution, (c) size distribution of the Ag NC's and (d) absorption spectra of the samples .....	71

<b>Figure 3.8.</b> Schematic of the Ag NC's and Ag NP's obtention proposed by Hssen et al 2017. Size-dependence of Ag species in fluorophosphate glasses are driven by concentration, heat treatment and annealing time.....	72
<b>Figure 3.9.</b> (a) Schematic illustration of femtosecond laser writing in a transparent host matrix. When a high enough energy pulses is focused into a material, optical breakdown is observed. In this case, silver ions ( $Ag^+$ ) is reduced into silver nanoclusters ( $Ag_m^{x+}$ ); (b) Multiphoton absorption promotes electrons from the valence band to the conduction band; (c) optical microscope images (excitation at 375 nm) of linear Ag NC's formed after femtosecond laser irradiation. Adapted from the references [GATASS; MAZUR, 2008; SMETANINA, 2016].....	74
<b>Figure 4.1.</b> Daylight pictures (a), luminescence images excited by a UV lamp at 254 nm (b) and 365 nm (c) for all the glasses.....	81
<b>Figure 4.2.</b> (a) UV-visible and near infrared transmission spectra of glass samples as a function of $AgNO_3$ concentration with their corresponding photograph presented in inset. (b) Absorption spectra of Ag NCs doped prepared glasses.....	82
<b>Figure 4.3.</b> TEM images of Ag5 (a and b) and Ag10 (c and d) samples as well as corresponding diameter distribution (e) and (f), respectively. Selected area electron diffraction pattern (SAED) taken from Ag5 and Ag10 samples.....	84
<b>Figure 4.4.</b> Normalized excitation spectra (PLE). The emission was fixed at 550 nm.....	85
<b>Figure 4.5.</b> Normalized emission spectra of the samples excited at 360 nm; (b) chromaticity diagram.....	86
<b>Figure 4.6.</b> Normalized (a) excitation and (b) emission spectra of the fluorophosphate glass doped with 7mol% $AgNO_3$ monitored at different wavelengths, respectively.....	88
<b>Figure 4.7.</b> Normalized emission spectra (a) the CIE chromaticity diagram (b) and (c) luminescence images showing multicolor emission of 5 mol% $AgNO_3$ doped sample excited/monitored at different wavelengths.....	90
<b>Figure 4.8.</b> Luminescence decay curves for different detected wavelengths upon excitation at 360 nm.....	92
<b>Figure 4.9.</b> Schematic configuration coordinate diagram of $Ag_m^{x+}$ NC's.....	94
<b>Figure 5.1.</b> Absorption spectra of the pure (black line), and co-doped with $Pr^{3+}/Yb^{3+}$ and different Ag amount. Increasing the Ag amount, the Ag NC's absorption increases and overlaps the $Pr^{3+}$ absorption between 400 and 500 nm. The inset shows $Pr^{3+}$ absorption lines in the visible range for the sample PYAg0 and the digital pictures show the change in color as a function of the Ag content.....	101

<b>Figure 5.2.</b> TEM images of the samples PYAg2 (a), PYAg4 (d), PYAg10 (g) and selected area electron diffraction (b), (e) and (h) for the respective samples (right side). Figures (c), (f) and (i) represent the size distribution of the NC's and NP's.....	104
<b>Figure 5.3.</b> Excitation spectra with emission wavelength fixed at 564 nm (a); excitation with emission fixed at 606 nm ( $\text{Pr}^{3+}$ lines) and 980 nm ( $\text{Yb}^{3+}$ lines), (b) and (c), respectively. Deconvolutions of the excitation bands for the samples PYAg2 (d), PYAg4 (e) and PYAg10 (f).....	106
<b>Figure 5.4.</b> Emission spectra in visible range (a–d) and near-infrared (e–h) performed for excitation in silver species absorption. The broadband emissions in the visible range characterize the emissions of Ag and $\text{Ag}_m^{\text{nt}}$ species and an ET is observed by the $\text{Pr}^{3+}$ emission at 500 and 606 nm (figure 5b-d). $\text{Yb}^{3+}$ emission resulting from ET, observed at different Ag NC's excitation (e-h).....	108
<b>Figure 5.5.</b> (a) Visible emission of co-doped ( $\text{Pr}^{3+}/\text{Yb}^{3+}$ ) and different Ag amount in fluorophosphate glasses; (b) time resolved spectra for the sample PYAg10.....	110
<b>Figure 5.6.</b> Decay curves of $\text{Pr}^{3+}$ ions with emission fixed at 486 nm ( $^3P_0 \rightarrow ^3H_4$ ) and 606 nm ( $^1D_2 \rightarrow ^3H_4$ ) under excitation at 444 nm, figures (c) and (d), respectively. The red lines are double exponential fit.....	112
<b>Figure 5.7.</b> (a) Near-infrared emission obtained for the sample without Ag (PYAg0) and with different Ag amount upon excitation at 444 nm; (b) experimental decay curves of the samples upon excitation at 444 nm and emission at 980 nm. The red line are double exponential fit.....	113
<b>Figure 5.8.</b> Partial diagram of energy levels of co-doped $\text{Pr}^{3+}/\text{Yb}^{3+}$ and Ag nanoclusters in fluorophosphate glass.....	114
<b>Figure 6.1.</b> Absorption spectrum of the sample Ag5.....	121
<b>Figure 6.2.</b> DSC curve of Ag5. The $T_g$ is indicated by the intersection point between the two red lines at 421 °C. The onset of crystallization is at 483 °C and the maximum crystallization occurs at 512 °C.....	122
<b>Figure 6.3.</b> Transmission electron microscopy (TEM) of the sample Ag5. Black spots characterize Ag NC's, whose show absence of cristalinity according to selected area electron diffraction (SAED), figure (b). Size distribution of the Ag NC's was obtained by analysis of 250 nanoclusters, figure (c).....	123
<b>Figure 6.4.</b> (a) PLE of the sample collected at different emission; (b) PL at different excitation wavelength and chromaticity diagram (c).....	124
<b>Figure 6.5.</b> (a) Partial setup for direct laser writing (DLW). The laser was perpendicularly focused by an objective on the top level of the sample and the patterns were written by displacements on the xy plane; (b) digital picture from the sample under UV light exposure; (c) fluorescence under UV light of all patters obtained after DLW. All patterns have the same dimensions 100 x 100 $\mu\text{m/s}$ , and the horizontal and vertical	

distance among them is around 100  $\mu\text{m/s}$ ; (d) expanded images for the patterns recorded at 10 $\mu\text{m/s}$  at different voltages. All the images show double line of Ag nanocluster under UV irradiation in dark field; (e) integrated fluorescence intensity ranging from 0,6 to 0,9 V.....126

**Figure 6.6.** (a) Transmission spectra recorded after different fs laser dose and (b) microfluorescence for the same parameters (except for the pristine zone, not recorded).....127

**Figure 6.7.** (a) Partial configuration for WG writing. The y axis may be taken as reference for the laser propagation; (b) depth of the WG (side view); (c) double line of Ag NC's after laser passing at 300  $\mu\text{m/s}$ ; (d) refractive index profile after formation of Ag NC's double line and (e) mode profile of the WG under 630 nm laser input.....128

## LIST OF TABLES

<b>Table 1.1</b> Number ( $n$ ) of electrons distributed in the unfilled $4f$ orbital of $\text{Ln}^{3+}$ ions and the electronic configuration. Adapted from the reference.....	<b>31</b>
<b>Table 2.1</b> Experimental decay time values for the samples under excitation at 521 nm.....	<b>56</b>
<b>Table 4.1</b> Experimental decay time values of 5 mol% $\text{AgNO}_3$ doped sample at different detected wavelengths.....	<b>92</b>
<b>Table 5.1</b> Molar composition of the raw materials and sample nomenclature.....	<b>100</b>
<b>Table 5.2</b> Experimental decay times of the samples under excitation at 444 nm and emission at 486 and 606 nm. The average time of the fast ( $\tau_1$ ) and slow ( $\tau_2$ ) decay is given by ( $\tau_M$ ).....	<b>110</b>

## LIST OF ABBREVIATIONS

Ag – Silver	HRTEM – High Resolution Transmission Electron Microscopy
Ag NC's – Silver Nanoclusters	kV – Kilovolt
Ag NP's – Silver Nanoparticles	LED – Light Emitting Diode
a-Si - Amorphous silicon	LLSC – Luminescent Layer for Solar Cell
CB – Conduction Band	Ln <sup>3+</sup> - Lanthanides
CET – Cooperative Energy Transfer	mc-Si - Multicrystalline silicon
CIE - Commission Internationale d'Eclairage	PA – Photon Avalanche
CR – Cross Relaxation	PV – Photovoltaic
CTB – Charge Transfer Band	SAED – Selected Area Electron Diffraction
DC – Downconversion	SC – Solar Cell
DLW – Direct Laser Writing	Si SC – Silicon Solar Cell
E <sub>g</sub> – Bandgap	Single-crystalline silicon – sc-Si
EMU – Energy Migration Upconversion	SPR – Surface Plasmon Resonance
ESA – Excited State Absorption	TEM – Transmission Electron Microscopy
ET – Energy Transfer	T <sub>g</sub> – Transition Temperature
ETE – Energy Transfer Efficiency	UC – Upconversion
ETU – Energy Transfer Upconversion	UV – Ultraviolet
FWHM – Full Width of Half-Maximum	Vis – Visible
GSA – Ground State Absorption	WG – Waveguide

## TABLE OF CONTENTS

<b>Chapter 1. Luminescent layers for solar cell application.....</b>	<b>24</b>
1 Introduction.....	24
1.2 Lanthanides.....	30
1.3 Energy levels of RE <sup>3+</sup> .....	32
1.4 Upconversion.....	34
1.5 UC for photovoltaic application.....	37
1.6 Materials for upconversion.....	42
1.7 References.....	43
<b>Chapter 2. UC mechanisms under 1550 nm excitation in Er<sup>3+</sup>- doped fluoroindate glasses.....</b>	<b>48</b>
2.1 Introduction.....	48
2.2 Experimental section.....	50
2.3 Results and discussion.....	51
2.3.1 Photoluminescent properties.....	51
2.3.2 Photocurrent analysis.....	57
2.4 Conclusions.....	58
2.5 References.....	59
<b>Chapter 3. A brief introduction to the silver nanoclusters.....</b>	<b>62</b>
3.1 Silver nanoclusters.....	62
3.2 Materials for Ag NC's stabilization.....	67
3.3 Effect of temperature, concentration and femtosecond laser radiation on the Ag NC's stabilization.....	68
3.3.1 Temperature dependence.....	68
3.3.2 Concentration dependence.....	70
3.3.3 Ag NC's induced by laser radiation.....	72
3.4 References.....	75

<b>Chapter 4. White light generation in fluorophosphate glasses doped Ag nanoclusters.....</b>	<b>78</b>
4.1 Introduction.....	78
4.2 Experimental section.....	79
4.3 Results and discussion.....	80
4.3.1 Optical properties and nanostructure analysis.....	80
4.3.2 Excitation and emission spectra dependences on Ag nanoclusters concentrations.....	85
4.3.3 Emission spectra dependence on excitation wavelengths.....	89
4.4 Conclusions.....	94
4.5 References.....	95
<b>Chapter 5. Photoluminescence of silver nanoclusters in co-doped Pr<sup>3+</sup>,Yb<sup>3+</sup> fluorophosphate glasses.....</b>	<b>98</b>
5.1 Introduction.....	98
5.2 Experimental section.....	100
5.2.1 Glass preparation.....	100
5.2.2 Characterizations.....	100
5.3 Results and discussion.....	101
5.3.1 Absorption spectra.....	101
5.3.2 Transmission electron microscopy.....	102
5.3.3 Spectroscopy results.....	105
5.4 Conclusions.....	114
5.5 References.....	115
<b>Chapter 6. Microstructures tailored by direct laser writing (DLW) technique in Ag containing fluorophosphate glass: application in 3D waveguide.....</b>	<b>118</b>
6.1 Introduction.....	118
6.2 Experimental section.....	119
6.2.1 Glass synthesis and spectroscopy characterization.....	119
6.2.2 DLW structures and waveguide fabrication.....	120
6.3 Results and discussion.....	121
6.3.1 Structural characterization of the Ag NC's and optical properties.....	121

6.3.2 Micro-patterning of Ag NC's by DLW and optical characterization.....	124
6.3.3 Preparation and optical characterization of waveguide.....	129
6.4 Conclusions.....	130
6.5 References.....	131
<b>Chapter 7. General conclusions.....</b>	<b>132</b>
<b>Appendix I.....</b>	<b>133</b>

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## Luminescent layers for solar cell application

*During the past 30 years there has been a remarkable growth in the use of luminescent materials for solar cell application. A huge variety of amorphous and crystalline materials has been employed in order to increase the photocurrent of different photovoltaic systems. In this sense, lanthanides ions have been employed as good candidates for luminescent layers, which may be used in the rear and the top face of certain solar cells. In this chapter, we present an overview of the solar cells landscape, the drawbacks associated with the solar spectrum mismatch and the pathways to overcome such issues. A greater emphasis is given to the solar cell from the first generation that are silicon-based solar cells. Along this chapter, it may be note that this kind of cell is taken as example to explain an increment of efficiency using luminescent layers. However, it is far beyond of this chapter to deal with efficiency using an AM 1.5G conditions as well as the measure of short-circuit current, fill factor and other parameters related. For this purpose, whenever necessary, the reader is direct to some references throughout the chapter. Besides, rare-earth spectroscopy is also mentioned as well as one of the most explored properties, the “upconversion” phenomenon.*

### 1 Introduction

The large-scale use of fossil fuels has brought serious issues of global warming and environmental pollution due to the nonrenewable energy consumption [KANG et al, 2010; LUNT; BULOVIC, 2011]. Since the introduction of first commercial solar cell (SC) in the 70ths years, research efforts in SC scaled up the knowledge to sustainable sources. Presently, solar energy is an emerging and alternate source of energy that can direct the usage to the clean and abundant form of energy [AHMED; HABIB; JAVAID, 2015]. Moreover, solar energy is one of the few renewable, low-carbon resource with maturity to meet ever-growing global demand for electricity [JEAN et al., 2015].

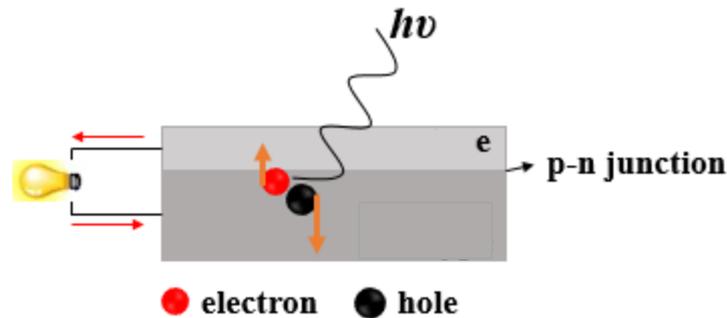
Daily, the sun delivers about 174 000 TW of energy to the upper level of the atmosphere, meaning 10 000 times more energy than our global energy consumption [RUHLE; SHALOM; ZABAN, 2011]. This energy corresponds to an average energy power 1366 TW.m<sup>-2</sup>. Taking into account atmosphere absorption and scattering, this value is reduced to 1000 W.m<sup>-2</sup>. Moreover, with further losses due to seasonal

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variation, oblique incident radiation, weather attenuation and cloud cover, leads to a global daily average of  $183 \text{ W.m}^{-2}$  [JACOBSON; DELUCCHI, 2011].

The most direct way to convert solar radiation into electricity is the use of photovoltaic (PV) cells, also commonly named as solar cells (SC). In such device, a semiconductor composed by a positive and negative (p-n) junction develops photovoltage. The detailed process to generate current is ascribed as following: when solar radiation (photons) is absorbed, electrons raise to a higher energy state, and then they flow to an external circuit, since electron-hole recombination pair is avoided. Electron and holes move to opposite electrodes and their recombination generate electric current. Such photovoltaic effect happens when few conditions are observed: (i) photons absorbed by the solar cell must to have energy equal or higher than the bandgap ( $E_g$ ) of the cell; (ii) the absorption of photons must to guarantee the formation of electron-hole pair, also called charge carriers; (iii) spatial separation of the charge carriers is observed by a p-n junction, so that any recombination is prevented; (iv) after electron flowing in the external circuit, electron-hole must to be regenerated again on the positive layer. For the sake of brevity, only main factors to generate photocurrent are mentioned above, however, the electric current density is affected by factors like shading, cell temperature, cable thickness, solar radiation intensity, climate conditions, cell lifetime and other [CHOPRA; PAULSON; DUTTA, 2004; GOETZBERGER; HEBLING; SCHOCK, 2002]. Moreover, Figure 1.1 represents the simplest configuration in which the SC is based in a double-layer (p-n junction) diode. Along this chapter, we are going to show complexes system, like those shown in Figure 1.2, in which several layer may be present in a PV device.

Figure 1.1. Schematic of a single-junction solar cell.



Along the last 5 decades, solar cell technology has grown dramatically owing the evolution of new materials that found their specific niche through different SC generation. In 2015, Jean et al. has presented a review paper regarding solar cell technology where the classification follows a trend in which the primary light-absorbing material was taken into account. According to them, two categories are defined: wafer-based cells and thin-film cells, respectively. Still, the generations are also referred as first, second and third generation, where the first one refers to the wafer-based SC, the second one refers to thin-film SC and the third one refers to the commercial thin-film SC. Here, we adopted the same definition described by Jean et al. [JEAN et al, 2015].

Figure 1.2 a-c shows the SC technology according wafer-based and thin-film classification. Wafer-based SC, or also called first generation, comprises the most mature of all photovoltaic technology and commercially speaking, that is silicon SC (Figure 1.2a). Silicon SC are classified as single-crystalline (sc-Si), multicrystalline (mc-Si) and amorphous (a-Si). Record cell efficiency for sc-Si, mc-Si and a-Si stand at 25.6%, 20.8% and 10.2% [GREEN, 2015], respectively. Nowadays, the global production counts with 90% of c-Si. In a general view, Si SC is composed by a positive (p-Si) and negative layer (n-Si) in which the thickness of the negative is smaller than the positive one. Back and frontal contacts are constituted of aluminum (Al) layer and thin strips of deposited silver (Ag). Other two SC compose the first generation, they are GaAs (Gallium arsenide) and III-IV MJ (Multijunction), Figure 1.2a. The major feature of III-IV MJ is the multiphoton absorption. In other words, this SC is able to absorb in different bandgaps across the solar spectrum. Two or more single-junction

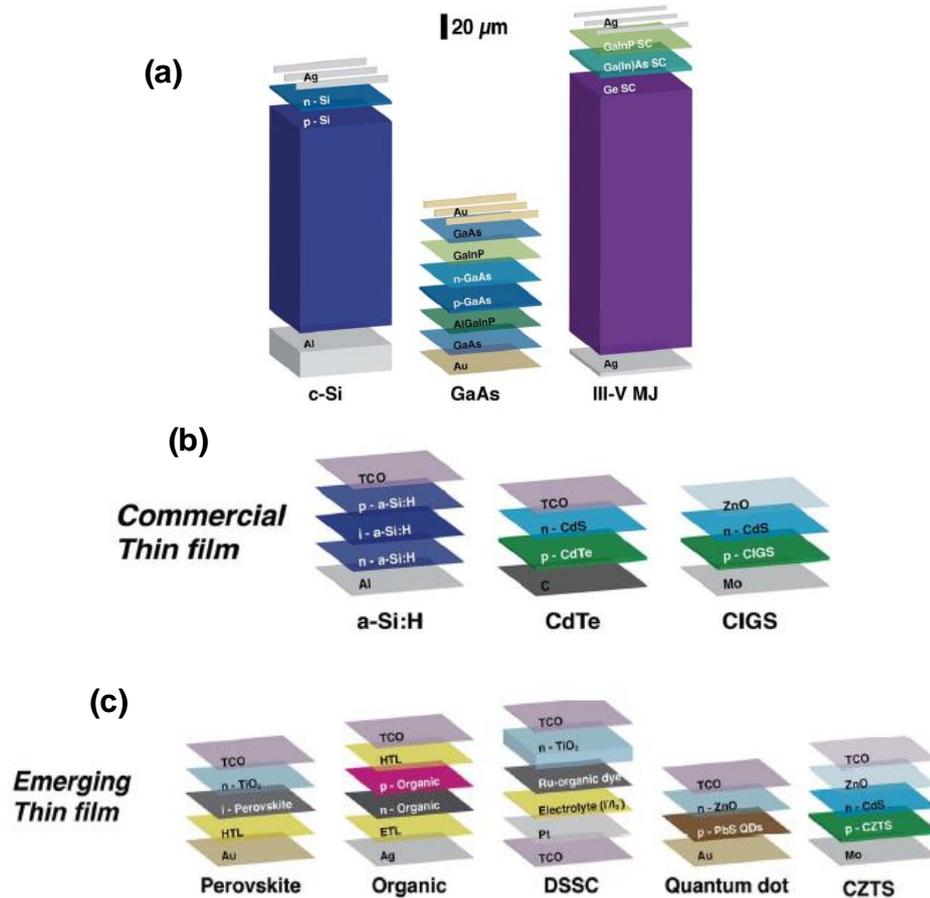
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layer composes this cell, in which each layer is stacked by decreasing the energy gap of the semiconductor. The main advantage of this configuration is the reduction of thermalization losses and the efficiency for two, three and four junctions can reach 46%, 44.4% and 34.1%, respectively. However, high cost materials and complex manufacturing limit a high-scale production.

Figure 1.2b shows the second SC generation, which is composed by commercial thin-film like a:Si-H (Hydrogenated amorphous silicon), CdTe (Cadmium telluride) and CIGS (Copper indium gallium diselenide). Cell efficiency for these SC are 12.2% for a:Si-H, 21% for CdTe and 21.7% for CIGS. A higher absorption efficiency is the key advantage compared to Si SC, and this allows the use of thin films of few microns of thickness. Besides, low material application may reduce the final cost; also, transport and installation may become easy task for this technology. However, limiting factor like high toxicity of cadmium, telluride; low abundance of indium and light-induced degradation for a:Si-H can hinder the large-scale application of these cells [JEAN et al, 2015].

Third generation extends the concept of thin film for nanostructured materials such as organic, perovskite, quantum dots, dye-sensitized solar cell (DSSC) and copper zinc tin sulfide ( $\text{Cu}_2\text{ZnSnS}_4$ , CZTS), Figure 1.2c. Key advantages for this generation is the ability to form flexible, light and thin SC, simple process manufacturing, low-temperature processing compared to those SC from the first generation and second [HAMAN et al, 2008; GREEN et al, 2002]. Solar efficiency for these SC are 20,1% (Perovskite), 11,1% (Organic), 12,3% (DSSC), 9,2% (Quantum dots) and 12,6% (CZTS). Despite the high plenty of organic elements that is advantageous for large-scale deployment, these SC, naturally face constraints issues like low durability and toxicity in Perovskite cells, low lifetime for all SC compared to the Si SC, in which each module has medium lifetime of 25 years [QIAO et al, 2018]. Still, unlike Si, cells from the second and third generation suffer with scarce elements and photodegradation under long time ultraviolet dose.

Figure 1.2. Solar cell generation according to wafer-based and thin-film classification. (a) First generation: wafer-based; (b) second generation: commercial thin film; and (c) second generation: thin-film.

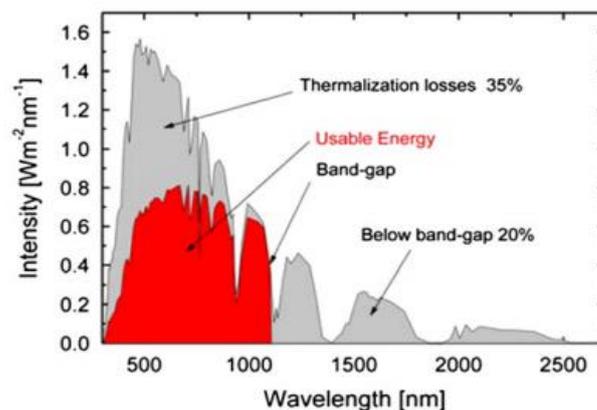


SC technology landscape still drives attention and efforts to cells from the first generation, especially, Si SC. Silicon counts with high abundance in earth crust, non-toxicity and higher cell efficiency. However, one major drawback hinder the enhancement of the efficiency that is solar spectrum mismatch. The maximum theoretic efficiency demonstrated for a single junction SC is 32%, this limit was first theoretically established by Shockley-Queisser [SHOCKLEY; QUEISSER, 1960]. Indeed, SC efficiency is limited also by the fact that sunlight reaches the surface cell in a non-concentrated and diffuse form, so that external devices (mirrors) are required to focus the radiation direct to the cell [CORREIA et al., 2014].

Solar spectrum mismatch brings two consequences to SC with a specific bandgap ( $E_g$ ), that is the thermalization of charge carriers and sub-bandgap loss, or also called loss by transmitted photons. Thermalization process is generated by absorption of high-energy photons, photons with energy higher than SC  $E_g$ . Meanwhile, sub-bandgap losses is generated by transmission of lower energy-photons, photons with lower energy than SC  $E_g$ . Both kind of photons, higher and lower energy than SC  $E_g$ , cannot be efficiently converted in electric energy. Taking as example Si SC with a bandgap of 1.1 eV (1100 nm), a broad part of the solar spectrum will be lost by thermalization and transmission. Solar spectrum on the surface earth widely ranges from 300 to 2500 nm with maximum intensity around at 500 to 600 nm [RODRÍGUEZ-RODRÍGUEZ, 2016] as it is shown in Figure 1.3. Grey color depicts the solar radiation incident in the earth's surface and the red region depicts the Si sensibility on the solar spectrum. Part of the incident energy (35%) is lost by thermalization and 20% by transmission (sub-bandgap photons) [ASIM et al., 2012].

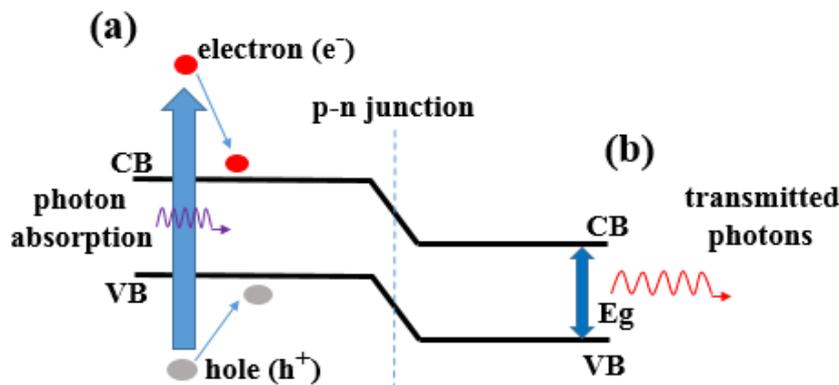
Losses process by thermalized photons and transmitted photons can be represented by the scheme in the Figure 1.4. Thermalization process are generated when charges arise straightforward to the edges of the conduction band (CB) and valence band (VB), process (a). Such excess of energy is readily dissipated as heat and its loss is substantial for SC with small bandgap. Contrarily, the SC cannot absorb transmitted photons, since its energy is not suitable to generate an electron-hole pair. The process (b) represents loss process by transmitted photons [RICHARDS, 2006].

Figure 1.3. Solar spectrum at earth's surface. The red region emphasizes the portion of the terrestrial sunlight that is absorbed by a silicon SC with bandgap of 1.12 eV, and spectral regions that are lost by thermalization and below bandgap or transmitted photons [ASIM et al, 2012].



Spectral mismatch in solar cell motivated many efforts to develop new materials with specific features to overcome such limitation [SARK et al, 2014]. The optimism in SC science inflates with the discover of the enhancement of efficiency conversion by modifying the input spectrum. Such conversion would be possible using lanthanide-based materials in which exhibit upconversion and downconversion phenomenon [CORREIA et al, 2014; ENDE; AARTS; MEIJERINK, 2009]. The next topic, we present a brief physics and chemistry introduction concerning lanthanides elements as well as their application in solar cell. For the sake of brevity, only the most employed lanthanide elements will be treat. Still, we are going to restrict the application of lanthanides elements in amorphous materials as well as their use as luminescent layers.

Figure 1.4. Losses in a single-junction SC. Process (a), thermalization photons and (b), transmitted photons. Adapted from reference [RICHARDS, 2006].



## 1.2 Lanthanides

Lanthanides ions ( $\text{Ln}^{3+}$ ) are commonly formed by ionization of an outer electron of atoms located at the bottom of the periodic table after lanthanum: from Cerium (atomic number 58) to Ytterbium (atomic number 70), which possess the electronic configuration  $5s^25p^65d^14f^16s^2$  and  $5s^25p^64f^{14}6s^2$ , respectively [SOLÉ; BAUSÁ; JAQUE, 2005]. Lanthanides ions have general configuration  $[\text{Xe}]4f^{n-1}5d6s^2$ , where [Xe] stands for the Xenon electronic structure and  $n$  is the number of electrons in the unfilled  $4f$  shell, that varies from 1 ( $\text{Ce}^{3+}$ ) to 13 ( $\text{Yb}^{3+}$ ). Although some ions may assume divalent ( $\text{Sm}^{2+}$ ,  $\text{Eu}^{2+}$ ,  $\text{Yb}^{2+}$ ) and tetravalent state of oxidation ( $\text{Ce}^{4+}$ ,  $\text{Tb}^{4+}$ ,  $\text{Pr}^{4+}$ ), the trivalent state is the most stable. To reach such ionization, the  $\text{Ln}^{3+}$  may lose a  $5d$ ,  $6s$

and one of the 4*f* electrons, leaving a partially filled 4*f* orbital that is shielded by the 5*s* and 5*p* orbitals [DWIVEDI; ZILIO, 2013]. Table 1 summarizes the lanthanide atoms in a trivalent configuration, the number of electrons valence on the 4*f* orbital of each atom and the electronic configuration as well.

Table 1.1: Number (*n*) of electrons distributed in the unfilled 4*f* orbital of Ln<sup>3+</sup> ions and the electronic configuration. Adapted from the reference [SOLÉ; BAUSA; JAQUE, 2005].

Atomic Number	(Ln <sup>3+</sup> ) ion	<i>n</i>	Electronic configuration (Ln <sup>3+</sup> )
58	Ce <sup>3+</sup>	1	4 <i>f</i> <sup>1</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
59	Pr <sup>3+</sup>	2	4 <i>f</i> <sup>2</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
60	Nd <sup>3+</sup>	3	4 <i>f</i> <sup>3</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
61	Pm <sup>3+</sup>	4	4 <i>f</i> <sup>4</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
62	Sm <sup>3+</sup>	5	4 <i>f</i> <sup>5</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
63	Eu <sup>3+</sup>	6	4 <i>f</i> <sup>6</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
64	Gd <sup>3+</sup>	7	4 <i>f</i> <sup>7</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
65	Tb <sup>3+</sup>	8	4 <i>f</i> <sup>8</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
66	Dy <sup>3+</sup>	9	4 <i>f</i> <sup>9</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
67	Ho <sup>3+</sup>	10	4 <i>f</i> <sup>10</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
68	Er <sup>3+</sup>	11	4 <i>f</i> <sup>11</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
69	Tm <sup>3+</sup>	12	4 <i>f</i> <sup>12</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>
70	Yb <sup>3+</sup>	13	4 <i>f</i> <sup>13</sup> 5 <i>s</i> <sup>2</sup> 5 <i>p</i> <sup>6</sup>

Since the orbitals 5*s*<sup>2</sup> and 5*p*<sup>6</sup> are totally filled electronically, optical and magnetic properties are ruled by the unfilled 4*f* shell and in some ions like Ce<sup>4+</sup> and Eu<sup>2+</sup>, by the 5*d* vacant orbital [REID, 2016]. However, yttrium (Y) also has +3 oxidation state and does not show optical transitions. In addition, Lanthanum (La<sup>3+</sup>) and Lutetium ions (Lu<sup>3+</sup>) have a completely empty and a completely filled 4*f* shell, therefore they are not optically active. Moreover, the 4*f*-4*f* direct interaction between neighboring rare-earth

ions may be considered as negligible and the behavior of the 4f electrons seems to be different from that s, p and d electrons [WYBOURNE, 2004].

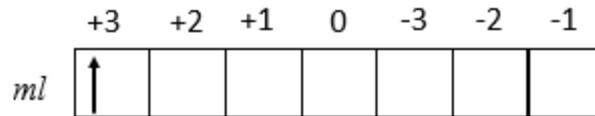
The shielded effect of the 4f orbital results in a poor interaction with the closest neighbor in a crystalline medium, by consequence, these inner electrons do not participate directly in chemical bonds. Such case, leads to the well-known situation “*weak crystalline field*”, which an ion is slightly perturbed by the crystalline field. The effect of the crystal field is to produce a slight shift in the energy levels and to cause additional level splitting within the *f* shell.

The electrons confined within the 4f orbital experiment a rearrange that is parity-forbidden. However, owing to a forced dipole electric transitions, most of the radiative transition become partially allowed by mixing the 4f orbitals with other orbitals that have different parity. In this sense, a promotion of a 4f electron into the 5d orbital are allowed by parity rule. In addition, a charge transfer band (CTB) can also implies an electronic transition in a trivalent lanthanide. This transition is allowed by the Laporte’s selection rule [BUNZLI, 2007].

Taking into account the low perturbation of the 4f electrons by a electric field, the *f-f* transitions become quite narrow and sharp compared to metal ion environment that directly affects the electrons involved on the optical properties. Consequently, a RE<sup>3+</sup> ion may have similar optical spectra among several different crystalline or amorphous medium [JORGENSEN; REISFELD, 2005; TANABE, 2015].

### 1.3 Energy levels of the RE<sup>3+</sup>

When a RE<sup>3+</sup> ion is dispersed in a crystalline host, it will experiment an electric field that will slightly perturb the  $^{2S+1}L_J$  states, where S and L are the total spin and total orbital quantum number, *J* is the total angular quantum number, which can assume the values  $J = (L+S); (L+S) - 1; (L+S) - 2; \dots; (L-S)$  [SOLÉ; BAUSÁ; JAQUE, 2005]. In spectroscopic notation, *L* may assume different symbols that are derived from the spectroscopic lines corresponding to s, p, d, f (sharp, principal, diffuse and fundamental) orbitals. For example, we may predict the ground state to the Ce<sup>3+</sup> ion, which has an electronic configuration  $4f^1 5s^2 5p^6$ . According to the Pauli Exclusion Principle, this only one electron 4f may be distributed in the following diagram box:

Figure 1.5. Electronic distribution of the  $Ce^{3+}$  ion in a box diagram.

Since  $ml = +3$  and  $L = \sum ml$ , the orbital quantum number ( $L$ ) will be 3. Each  $L$  is associated to a symbol that is characteristic of the spectroscopic lines corresponding to the orbitals sharp, principal, diffuse and fundamental ( $s, p, d, f$ ).

$$L = \begin{matrix} 0 & 1 & 2 & 3 & 4 & 5 \\ S & P & D & F & G & H \end{matrix} \quad (1)$$

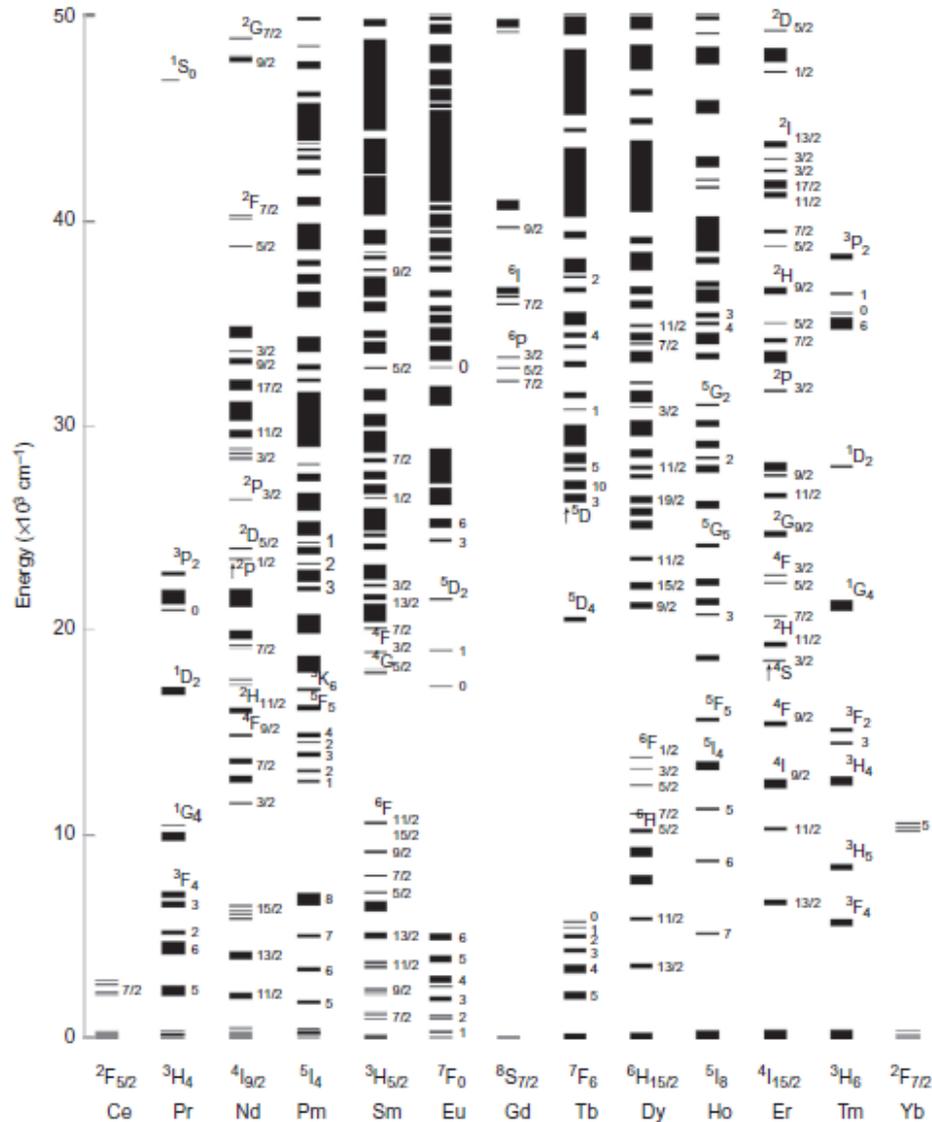
Assuming that  $L = 3$ , the electron from the  $Ce^{3+}$  will be in an orbital associated to a quantum number  $F$ . Another quantum number to determine is spin total ( $ms$ ) that is associated to the electronic orientation and can have positive and negative values. Assuming a positive configuration to the spin electronic ( $+1/2$ ), the  $ms$  will be equal to  $+1/2$ , and considering that  $S = \sum ms$ , the total spin is  $S = +1/2$  and  $2S+1$  is equal to 2. As  $J = (L+S); (L+S) - 1; (L+S) - 2; \dots; (L-S)$  and  $S$  is  $+1/2$ , the state with the lowest  $J$  value ( $L-S$ ) is  $5/2$ . Therefore, spectroscopic term for the ground state of the  $Ce^{3+}$  ions is  $^2F_{5/2}$ .

The barycenter of energy of each  $^{2S+1}L_J$  states may change according to the electric field surround the  $Ln^{3+}$ , since the electric field depend on the site symmetry and the nature of the ions in the crystal [YUHUA, 2015]. However, since the  $Re^{3+}$  ions assume a free-ion behavior, the position of the multiplets are almost independent of the crystal. Therefore, the gross feature of the energy level diagram remain unchanged [SOLÉ; BAUSÁ; JAQUE, 2005]. The maximum number of sublevels components for each  $^{2S+1}L_J$  level is given by  $2J+1$  when  $J$  is an integer and  $2J+1/2$ , when  $J$  is half-integer.

The absorption spectra of the  $Re^{3+}$  ions in lanthanum chloride crystal were target of extensive research by Dieke's group [BUNZLI, 2007]. The entire  $^{2S+1}L_J$  levels of energy of all the trivalent ions are shown in the Figure 1.6. The distribution varies from  $Ce^{3+}$  ( $4f^1$ ) to  $Yb^{3+}$  ( $4f^{13}$ ). The short horizontal lines indicates the energies of the  $^{2S+1}L_J$  multiplets and the width of the lines representing the crystal-field splittings of the multiplets. The largest split among the states by the order of  $10^4 \text{ cm}^{-1}$  are caused by

Coulomb interaction. The spin-orbit (S-L) interaction splits the terms in multiples separated by the order of  $10^3 \text{ cm}^{-1}$ . For the last, the crystal-field interaction split the multiplets by order of  $10^2 \text{ cm}^{-1}$  [WHEG et al, 2000].

Figure 1.6. Energy level for trivalent rare-earth ions ( $\text{RE}^{3+}$ ) in  $\text{LaF}_3$  crystal [ELISEEVA; BUNZLI, 2010].



## 1.4 Upconversion

One of the most explored spectroscopic properties of the  $\text{RE}^{3+}$  is the upconversion (UC). The origin of the first research in UC can be date in 1959 by Bloembergen, who proposed lanthanide-doped materials as infrared quantum counters [BLOEMBERGEN, 1959]. Nowadays, the interest in this phenomenon is far beyond

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and covers a wide research field. A huge variety of potential applications is explored, for example, bio-imaging [LIU; HOU; GAO, 2014], white light modulation [KARATAIRI; HOPPE, 2016; ELISEEVA; BUNZLI, 2011], thermal sensors [LALLA et al, 2015], photocatalysis [ULLAH et al, 2017], spectral converters for sola cell [GOLDSCHMIDT; FISCHER, 2015; SHPAISMAN, 2008], solid-state lasers and others [LIAN et al., 2013; -19; DONG, 2017; VASILIEV, 2016].

The term UC refers to a non-linear process in which the continuous absorption of two or more low-energy photons leads to the emission of high-energy photons (anti-Stokes emission) [REID, 2016]. Figure 1.7 summarizes the simplified mechanisms of UC emissions, which may be described as following: the first step corresponds in an excitation through a near-infrared source, indicated by the orange arrow in the Figure 1.7 (a)-(e). In this step, one electron that is in its ground state ( $E_0$ ) is excited and occupies the first metastable state ( $E_1$ ). This first step is fundamental and common to all the mechanisms; excited state absorption (ESA), energy transfer upconversion (ETU), photon avalanche (PA), cooperative energy transfer (CET) and energy migration upconversion (EMU). The second step differs from each mechanism. For example, the first mechanism (Figure 1.7a), the UC occurs in a single ion. After absorption of one photon, the same excited electron (now in  $E_1$  state) absorbs another photon of the same energy and goes to a further higher excited state (named as  $E_2$ ). From this state, radiative transition is observed with emission of a higher-energy radiation.

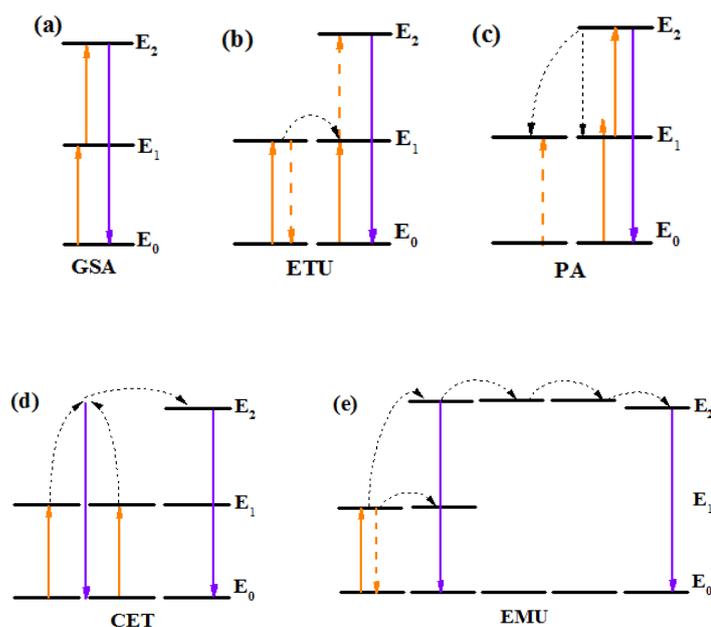
In an ETU mechanism (Figure 1.7b), two ions are involved and they are commonly named as sensitizer and activator, in which the sensitizer, generally, has a higher cross section than the activator. Since the first intermediated state from the sensitizer and activator are spaced in a close distance (resonant), the sensitizer transfer energy to adjacent ion that is excited to the  $E_2$  state. Commonly, the mechanism that happens between two ions is called “*upconversion cooperative system*” and can be observed between two similar ions or different ions [SUN, 2015]. However, for different ions the primordial condition is that the metastable state of the two ions must to be energetically close.

When the energy gap between an intermediated state and the ground state is non-resonant with the pump power, the PA can produce population in an excited state whose energy exceeds that the pump photon [JOUBERT et al, 1999; STRUMPEL et

al, 2007], as it is shown in the Figure 1.7c. This electron will be excited to a higher state energy ( $E_2$ ) by an ESA. Then, a cross relaxation (CR) excites an adjacent ion and the level of the same ion. It is peculiar to mention that a UC photon can be generated without a GSA step. In this case, the consumption of superexcited ions must to be less than the ground state ions [SUN, 2015].

Another mechanism that takes place in a cooperative sense is the CET, shown in the Figure 1.7d. In this process, two sensitizer ions are needed to sensitize the activator. This is due to an absence of a long-lived metastable state of the sensitizer. The last mechanism is the EMU, shown in the Figure 1.7e. In 2011, Liu and coworkers introduced this new UC mechanism based in core-shell structured nanoparticles. In the EMU, the energy transfer from the core to the shell follows a cascade-type sensitization in which the sensitizer is located on the core and the activator on the last shell [WANG, 2011].

Figure 1.7. Simplified energy levels for different two photons mechanisms in UC process. The excitation of the sensitizers ions are represented by orange arrows, non-radiative transition are represented by black dotted arrows and the purple arrows represent visible emission from the excited state  $E_2$  to the ground state  $E_0$ , respectively.

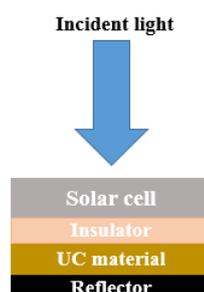


Despite the benefits resulting from the UC emissions (see applications in the beginning of this section), the main limitation of the  $RE^{3+}$  is the narrow and weak absorption associated with the nature of the  $4f-4f$  transitions [BUNZLI; PIGUET, 205], which is shielded by the orbitals  $6s$  and  $6p$ . For comparison, the  $Yb^{3+}$  that is the most efficient sensitizer used for generation of UC emission, has approximately a cross section of  $10^{-20} \text{ cm}^2$ , that is 1000 to 10000 lower than a dye molecule ( $10^{-17} - 10^{-16} \text{ cm}^2$ ). Still, the narrow spectral absorption, over 10 times narrower than an organic dye, compromises their utility for several applications. For example, the excitation using natural sources, like the sun, is not feasible to generate UC emission. Indeed, this phenomenon happens by using commercial sources, like diode laser operating at 808, 980 and 1500 nm [RODRÍGUEZ-RODRÍGUEZ, 2013]. Last but not the least, the total efficiency of an UC system can reach a maximum 50% since two photons is necessary to provide one photon.

### 1.5 Upconversion for photovoltaic application

Bringing back the concept of solar mismatch in solar cell, here considering the transmitted photons, first mentioned on the introduction of this chapter, we are able to scale-up one of the most interesting application of the upconversion phenomenon. Trupke et al. [TRUPKE; GREEN, WURFEL, 2002] performed the first theoretical approach investigating UC in photovoltaic devices. The authors report a theoretical system in which an upconverter layer is located in the rear level of a bifacial solar cell and the harvesting of sub-band-gap photons allowed them to find 63,2% of quantum efficiency for concentrated sunlight and 47,6% for nonconcentrated sunlight. Figure 1.8 summarizes the system used by Trupke et al.

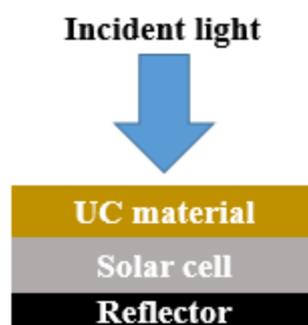
Figure 1.8. Scheme of an ideal system for UC in a bifacial solar cell proposed by Trupke et al. Adapted from the reference [TRUPKE; GREEN, WURFEL, 2002].



The packed system is composed by a bifacial solar cell, an insulator to isolate electronically the solar cell and the UC layer. Beneath the UC material, a reflector is located to reflect back the light emitted by the UC layer. Trupke et al. assumed that a photon selectivity would be applied to generate extra electron-hole pairs on the solar cell. In other words, the UC material would absorb the radiation transmitted by the first layer, here considered as solar cell. Then, the transmitted radiation absorbed by the UC material is delivered with a higher energy photon. The rear reflector located below the UC layer, reflects the light back to the solar cell. In this way, all the radiation converted by the UC layer is used to generate an electron-hole pair on the photovoltaic device.

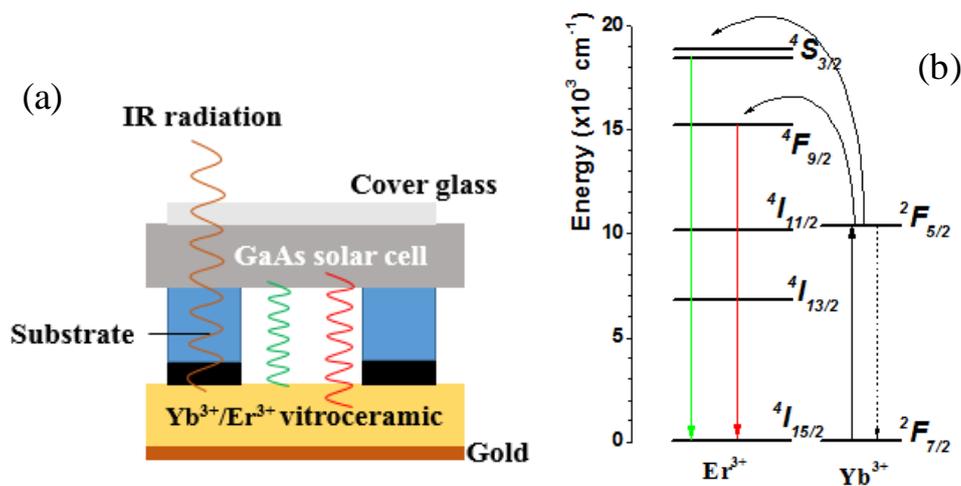
Viorel Badescu [BADESCU, 2008], in 2008, presented an extended study of the seminal work of Trupke et al. Badescu compared two different approaches: (i) a system composed by a solar cell and Upconversion layer (SC-UC) and (ii) upconversion layer and solar cell (UC-SC). In other words, the system (ii) presents the upconversion layer on the top level of the solar cell. Indeed, in this new approach, it is totally predictable that part of the radiation emitted by the UC layer, would be lost on the top surface (emission in all spatial directions). Moreover, part of the incident radiation would be lost by scattering on the UC layer. Therefore, no benefit can be expected when a UC layer is placed at the front of the solar cell.

Figure 1.9. Experimental scheme of an UC material and a solar cell proposed by Badescu. Adapted from the reference [BADESCU, 2008].



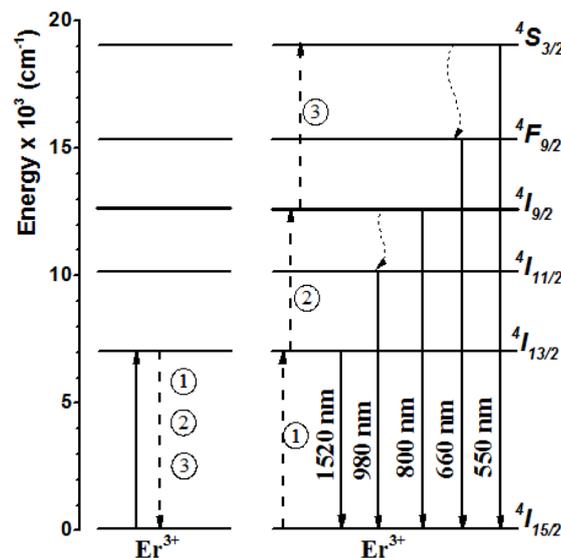
Few years earlier, in the 90<sup>th</sup> years, experimental insight has shown promising results by using the rare earth  $\text{Er}^{3+}$  and  $\text{Yb}^{3+}$  as upconverter materials in GaAs solar cell. In 1996, Gibart et al. [GILBART, 1996] developed a device based in a vitroceramic co-doped with  $\text{Yb}^{3+}/\text{Er}^{3+}$  to use part of the infrared radiation otherwise lost by transmission in a GaAs solar cell. Figure 1.10a shows the setup used and 1.10b shows the energy levels of the  $\text{Yb}^{3+}/\text{Er}^{3+}$  pair according to the mechanism proposed by the authors. An infrared laser (980 nm) irradiates the front face of a GaAs solar cell. The radiation pumping is called “*bellow bandgap IR irradiation*” because the radiation energy is lower than the bandgap of the solar cell, so that the incident radiation is transmitted to the upconverter material. By two-photon UC mechanism, the vitroceramic emits green and red radiations, according to the mechanism from the Figure 1.10b in which the  $\text{Yb}^{3+}$  ion absorbs two photons and after energy transfer to the closest  $\text{Er}^{3+}$ , then useful radiation is emitted to generates electron-hole pairs in a GaAs solar cell. In this experiment, a Ti-Sapphire IR laser at 1.39 eV lighted the GaAs solar cell, which has bandgap of 1.42 eV. For an input excitation of 1 W, the authors found an efficiency of 2.5% illuminating a substrate of 0.039 cm<sup>2</sup> of sensitive area

Figure 1.10. (a) Experimental setup of a GaAs solar cell coupled with  $\text{Yb}^{3+}/\text{Er}^{3+}$  co-doped vitroceramic, and (b) energy levels showing red and green emissions after energy transfer from  $\text{Yb}^{3+}$  to  $\text{Er}^{3+}$ . Figure (a) was adapted from the reference [GILBART, 1996].



Regarding to the most commercial deployed solar cell, silicon was first applied in a proof-of-principle in 2003, when Shalav et al. [SHALAV, 2007] reported a bifacial solar cell with the upconverter  $\text{NaYF}_4:\text{Er}^{3+}$  in the rear side of the device. Indeed, the rare earth  $\text{Er}^{3+}$  is one of the most useful ion for upconversion applications. The first excited state is the  $^4I_{13/2}$  and the energetic different between this state and the ground state ( $^4I_{15/2}$ ) is 0,827 eV that is equivalent to 1500 nm [CHEN, 2015]. Still, the states  $^4I_{9/2}$  and  $^4S_{3/2}$  are multiples of the  $^4I_{15/2} \rightarrow ^4I_{13/2}$  transition, so that an excitation at 1500 nm is ideally suited to populates the excited states and then releases visible and near-infrared photons due to its upconversion. Figure 1.11 shows UC between two ions  $\text{Er}^{3+}$  in which one is the sensitizer and the second one is the activator. Upon excitation around 1500 nm, the excited levels of the activator can be populated after subsequent energy transfer (ET) from the sensitizer, and then, radiative transitions can be observed through the emission of light at 980, 810, 660, 550 and 410 nm.

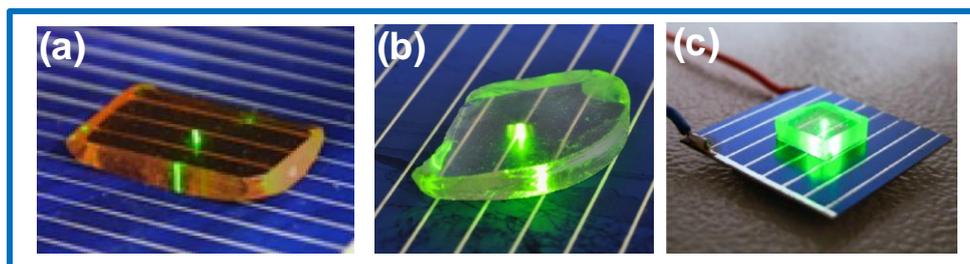
Figure 1.11. Energy levels of two  $\text{Er}^{3+}$  ions showing UC process upon excitation at 1500 nm. Three subsequent ET from the sensitizer to the activator guarantee the population of the excited states in which UC emission is observed.



Another proof-of-principle is that when a luminescent layer is placed on the top level of a solar cell. Despite proved by Badescu, that first mentioned no benefits can be achieved by placing a UC material in front of a solar cell, studies have shown an

increase in the photocurrent when the samples are excited between 1480 and 1550 nm [RODRÍGUEZ-RODRÍGUEZ et al, 2016; KRISHNAIAH et al, 2017; HENKE et al, 2009]. Indeed, in a commercial point of view, a device in which a layer present deleterious effect to the global efficiency is not desired, otherwise merely for research and developing new luminescent materials. Moreover, since the layer do not modify the solar cell features but only the incident radiation by spectral modification, applications of layers on the top level of solar cell are still attractive to stud photocurrent enhancement by UC properties. Figure 1.12 shows three experimental setup to measure photocurrent response when a UC layer is placed on the top level of a silicon solar cell. In Figure 1.12a, a tellurite glass doped with  $\text{Er}^{3+}$  and excited at 1500 nm; in 1.12b, a fluoroindate glass co-doped with  $\text{Yb}^{3+}$ - $\text{Er}^{3+}$  and is excited at 1480 nm and in 1.12c,  $\text{Er}^{3+}$  containing fluorozirconate glass is excited at 1540 nm. All the glasses are perpendicularly excited by a near-infrared source and the green color is due to the  $\text{Er}^{3+}$  transition ( $^4\text{S}_{3/2} \rightarrow ^4\text{I}_{15/2}$ ). In both cases, the samples are not use as encapsulated glass. In other words, samples are smaller than the active are of the solar cells, as we can see from the digital pictures (a), (b) and (c).

Figure 1.12. Working principle of a luminescent layer placed at the top level of a silicon solar cell. Adapted from the references [RODRÍGUEZ-RODRÍGUEZ et al, 2016; KRISHNAIAH et al, 2017; HENKE et al, 2009].

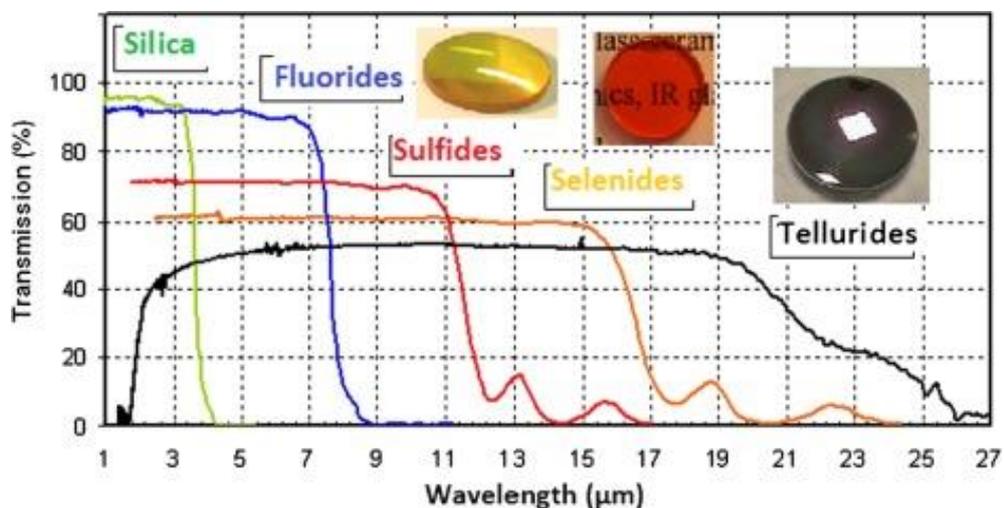


## 1.6 Materials for upconversion

There is a vast literature concerning monocrystalline and polycrystalline upconverter materials as luminescent layer for solar cell, in this section we introduced an overview regarding materials for potential upconversion application. Still, we restrict throughout the text to the use of the  $\text{Er}^{3+}$ -doped materials for UC in luminescent layers, as introduced in the last topics. As scribed before, lanthanide as  $\text{Er}^{3+}$ , for example, exhibit absorption in the near infrared. The main interest in materials for upconversion is that the window transparency extend to the visible and infrared domain [CALVEZ, 2017]. Visible transparency is defined by electronic transition from the valence band to the conduction band [NALIN et al, 2016]. In other words, the bandgap ( $E_g$ ), that is the energy gap between the valence and conduction band, which is associated to short wavelengths, defines the transparency in the visible range. On the other hand, vibrational modes defines the transparency in longer wavelength (infrared). Generally, massive atoms such as selenium (Se), sulfur (S) and tellurium extend the transmission band to longer wavelengths [CALVEZ, 2017].

Figure 1.13 presents the window transparency for some Chalcogenide glasses. The global transparency is governed by the electronic transition (short wavelength) and vibrational modes (longer wavelength). Silica and fluoride glasses have high transparency in the visible range (not shown), meanwhile sulfides, selenides and tellurides glasses have lower transparency in the visible range.

Figure 1.13. Window transparency for Chalcogenide glasses [CALVEZ, 2017].



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## General Conclusions

Fluoroindate and fluorophosphate glasses were prepared through the melting-quenching method. These two family of glasses were selected according to the physics and chemical properties. The low energy phonon matrix (fluoroindate) yielded a very intense upconversion emission for the sample with 7 mol% of  $\text{Er}^{3+}$ , in consequence, a high photocurrent was observed for low power excitation (22 mW). However, due to the low absorption and emission coefficient of the lanthanides, it is still far the application in spectral converters to increase solar cells efficiency. From the fluorophosphate glasses results, we conclude a size-dependence for the Ag NC's stabilization. The  $\text{F}^-$  ions played an important role on the Ag NC's reduction. In addition, the long molecular weight introduced by the metaphosphate groups may be a crucial reason for a non-agglomeration of Ag ions, since the phosphate groups may leave voids in the glass network. However, this is still in investigation in our research group. The dispersion of 5 mol% of Ag, seems to be an ideal concentration for Ag NC's obtention and SPR-free samples. This sample with 5 mol% showed a broadband emission that started in the UV region and extended until infrared at 730 nm. This broadband resulted from a size-distribution behavior in which the formation of nanoclusters responsible for the blue, green and red emission is observed. Still, silver nanoclusters showed an energy transfer (ET) to the rare-earths  $\text{Pr}^{3+}$  and  $\text{Yb}^{3+}$ . The composition with less amount of metaphosphate groups, allowed to disperse a less amount of silver, in consequence, 2 mol% was the optimum concentration to avoid the obtention of opaque glasses with SPR effect, as it was obtained for 4 and 10 mol%. However, the samples even for high amount present broadband emission in the visible range. For the last, very bright microstructures of silver were obtained after interaction of the sample with 5 mol% and the femtosecond laser. This sample showed to be a good candidate for waveguide application.