

Podolsky's electromagnetic theory on the null-plane gauge

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Abstract. Following the Dirac's technique for constrained systems we performed a detailed analysis of the constraint structure of Podolsky's electromagnetic theory on the null-plane coordinates. The null plane gauge condition was extended to second order theories and appropriate boundary conditions were imposed to guarantee the uniqueness of the inverse of the constraints matrix of the system. Finally, we determined the generalized Dirac brackets of the independent dynamical variables.

Keywords: Podolsky's electromagnetic theory, Null-Plane Gauge Conditions

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INTRODUCTION

The Podolsky's electromagnetic theory [1] was developed as a generalization of the Maxwell's electromagnetism. It adds to the Maxwell Lagrangian a quadratic term of the divergence of the field-strength tensor in order to avoid divergences such as the electron self-energy and the vacuum polarization current. The resulting theory is the only possible generalization of the electromagnetic field that preserves invariance under $U(1)$ [2]. The Lagrangian density for Podolsky's theory is

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - a^2\partial_\lambda F^{\mu\lambda}\partial^\gamma F_{\mu\gamma}, \quad F_{\mu\nu} \equiv \partial_\mu A_\nu - \partial_\nu A_\mu. \quad (1)$$

The theory yields field equations that are still linear in the fields, and predicts the existence of massive photons proportional to the inverse of the Podolsky's parameter a . The first consistent approach to the quantization of the field was made in instant-form by Galvão and Pimentel [3], meanwhile the authors studied the theory on the null-plane coordinates where Dirac's canonical formalism [4] was used to derived the correct gauge choice for this theory [5].

CANONICAL STRUCTURE

From the Lagrangian density (1), it is possible to calculate the canonical momenta

$$p^\mu = F^{\mu+} + 2a^2 \left(\eta^{\mu-} \partial_- \partial_\lambda F^{+\lambda} + \eta^{\mu i} \partial_i \partial_\lambda F^{+\lambda} - 2\partial_- \partial_\lambda F^{\mu\lambda} \right),$$

$$\pi^\mu = -2a^2\eta^{\mu+}\partial_\lambda F^{+\lambda}, \quad (2)$$

associated to the fields A_μ and $\bar{A}_\mu \equiv \partial_+ A_\mu$ respectively. From this relations we are able to define the primary constraints of the theory

$$\begin{aligned} \phi_1 &= \pi^+ \approx 0, \\ \phi_2^i &= \pi^i \approx 0, \\ \phi_3 &= p^+ - \partial_- \pi^- \approx 0, \\ \phi_4^i &= p^i - \partial_i \pi^- + F_{i-} - 4a^2 \partial_- [\partial_i \bar{A}_- - 2\partial_- \bar{A}_i + \partial_i \partial_- A_+ - \partial_j F_{ij}] \approx 0. \end{aligned} \quad (3)$$

The canonical Hamiltonian density can be expressed by

$$\begin{aligned} \mathcal{H}_c &= p^\mu \bar{A}_\mu + \pi^- (\partial_- \bar{A}_+ - \partial^i \bar{A}_i + \partial^i \partial_i A_+) - \frac{1}{2} (\bar{A}_- - \partial_- A_+)^2 - (\bar{A}_i - \partial_i A_+) F_{-i} \\ &\quad + \frac{1}{4} F_{ij} F^{ij} - a^2 (\partial_i \bar{A}_- - 2\partial_- \bar{A}_i + \partial_i \partial_- A_+ - \partial_j F_{ij})^2. \end{aligned} \quad (4)$$

With the canonical Hamiltonian $H_c = \int d^3x \mathcal{H}_c(x)$ and the primary constraints (3) we build the primary Hamiltonian

$$H_P \equiv H_c + \int d^3x u^a(x) \phi_a(x) \quad . \quad \{a\} = \{1, 2, 3, 4\}. \quad (5)$$

The consistency conditions on the constraints determine that the condition $\phi_1 \approx 0$ gives just the constraint $\phi_3 \approx 0$, which is already satisfied. The consistency for the remaining constraints give equations for some Lagrange multipliers and a secondary constraints appears: $\chi \equiv \partial_- p^- + \partial_i p^i \approx 0$. From this relation no more constraints can be found. The analysis leave us with the following set of constraints:

$$\begin{aligned} \chi &= \partial_- p^- + \partial_i p^i \approx 0, \\ \phi_1 &= \pi^+ \approx 0, \\ \phi_2^i &= \pi^i \approx 0, \\ \phi_3 &= p^+ - \partial_- \pi^- \approx 0, \\ \phi_4^i &= p^i - \partial_i \pi^- + F_{i-} - 4a^2 \partial_- [\partial_i \bar{A}_- - 2\partial_- \bar{A}_i + \partial_i \partial_- A_+ - \partial_j F_{ij}] \approx 0. \end{aligned} \quad (6)$$

It happens that χ and ϕ_1 are first-class constraints, while ϕ_2^i , ϕ_3 and ϕ_4^i are second-class ones. However, constructing the matrix of the second-class constraints we found that it is singular of rank four, which indicates that there must be a first-class constraint, associated with the zero mode of this matrix, and its construction is made from the corresponding eigenvector which gives a linear combination of second-class constraints. The combination happens to be just $\Sigma_2 \equiv \phi_3 - \partial_i \phi_2^i$ and it is independent of χ and ϕ_1 . Therefore, we have the renamed set of first-class constraints

$$\begin{aligned} \Sigma_1 &\equiv \pi^+ \approx 0, \\ \Sigma_2 &\equiv p^+ - \partial_- \pi^- - \partial_k \pi^k \approx 0, \\ \Sigma_3 &\equiv \partial_- p^- + \partial_i p^i \approx 0, \end{aligned} \quad (7)$$

and a set of irreducible second-class constraints

$$\begin{aligned}\Phi_1^i &\equiv \pi^i \approx 0, \\ \Phi_2^i &\equiv p^i - \partial_i \pi^- + F_{i-} - 4a^2 \partial_- [\partial_i \bar{A}_- - 2\partial_- \bar{A}_i + \partial_i \partial_- A_+ - \partial_j F_{ij}] \approx 0.\end{aligned}\quad (8)$$

The second-class constraints do not appear in the instant-form dynamics for this theory, thus, they are a common effect of the null-plane dynamics [7].

NULL-PLANE GAUGE FIXING AND DIRAC'S BRACKETS

Now let us eliminate the arbitrary functions from our theory by using gauge invariance to fix the remaining three degrees of freedom corresponding to the first class constraints (7). Let us choose the null-plane gauge defined by [8]

$$A_- \approx 0. \quad (9)$$

The time preservation of this condition is equivalent to setting

$$\partial_+ A_- = \bar{A}_- \approx 0, \quad (10)$$

whereas the consistency requires

$$\partial_+ \bar{A}_- = \partial_- \bar{A}_+ + \partial_i^x \bar{A}_i - \partial_i^x \partial_i^x A_+ \approx 0. \quad (11)$$

The relations (9, 10, 11) constitute the generalized null-plane gauge condition.

The next step is to calculate Dirac's Brackets which are defined by

$$\begin{aligned}\{A(x), B(y)\}^* &\equiv \{A(x), B(y)\} - \int d^3 z d^3 w \{A(x), \Theta_a(z)\} [(\mathcal{O}^{-1})]^{ab}(z, w) \\ &\quad \{\Theta_b(w), B(y)\}\end{aligned}\quad (12)$$

Here $\{\Theta_a(z)\}$ denote the set of ten second class constraints of the theory and $[(\mathcal{O}^{-1})]^{ab}(z, w)$ represent the inverse of the matrix formed with this set of constraints. The explicit evaluation of the inverse involves the knowledge of the inverse of the operators $(\partial_-^x)^{-1}$, $(\partial_-^x)^{-2}$, and $(\partial_-^x)^{-3}$. To achieve a unique solution it is necessary and sufficient to impose $\partial_-^x A_\mu = 0$, $\partial_-^x \partial_-^x A_\mu = 0$, and $\partial_-^x \partial_-^x \partial_-^x A_\mu = 0$ on $x^+ \rightarrow -\infty$ as the appropriate boundary conditions of the theory.

Thus, we have that the fundamental DB among the independent dynamical variables of the theory are:

$$\begin{aligned}\{A_+(x), \bar{A}_k(y)\}_D &= \frac{1}{32a^2} \partial_k^x \delta^2(x^\top - y^\top) \partial_-^x \int dv^- |x^- - v^-| |v^- - y^-| \\ \{A_+(x), \pi^-(y)\}_D &= -\frac{1}{2} \varepsilon(x^- - y^-) \delta^2(x^\top - y^\top) \\ \{A_k(x), p^l(y)\}_D &= \delta_k^l \delta^3(x - y)\end{aligned}$$

$$\begin{aligned}
\{\bar{A}_k(x), A_l(y)\}_D &= \frac{1}{16a^2} \delta_k^l |x^- - y^-| \delta^2(x^\top - y^\top) \\
\{\bar{A}_k(x), \bar{A}_l(y)\}_D &= -\frac{\delta_l^k}{128a^4} (x^- - y^-)^2 \varepsilon(x^- - y^-) (1 + 4a^2 \nabla_x^2) \delta^2(x^\top - y^\top) \\
&\quad + \frac{1}{16a^2} \partial_k^x \partial_l^x \delta^2(x^\top - y^\top) 2\partial_-^x \int dv^- |x^- - v^-| |v^- - y^-| \\
\{\bar{A}_k(x), p^l(y)\}_D &= -\frac{1}{16a^2} \varepsilon(x^- - y^-) \left[-\delta_k^l + 4a^2 \left(\partial_i^x \partial_k^x - \delta_k^l \nabla_x^2 \right) \right] \delta^2(x^\top - y^\top) \\
\{\bar{A}_k(x), \pi^-(y)\}_D &= -\frac{1}{4} \varepsilon(x^- - y^-) \partial_k^x \delta^2(x^\top - y^\top)
\end{aligned}$$

With these brackets we can deduce the fundamental ones that will lead, through the correspondence principle, to a consistent quantization of the field.

FINAL REMARKS

In the canonical analysis of the Podolsky's theory we found a set of three first-class constraints (7), followed by another set of four second-class ones (8) which are not present in the conventional instant-form dynamics [3]. The appearance of this kind of constraints is a common effect of the null-plane dynamics [7].

We eliminated the first class constraints (7) from our theory using the generalized null-plane gauge condition. By imposing the value of the field A_μ on the null-plane $x^+ = cte$, and the considered boundary conditions: $\partial_- A_\mu = 0$, $(\partial_-)^2 A_\mu = 0$, and $(\partial_-)^3 A_\mu = 0$ on $x^+ \rightarrow -\infty$, on $x^- = cte$, we have fixed the hidden subset of the first-class constraints [9] and got a unique set of Dirac's brackets among the fundamental fields of the theory.

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